

Gravity Flows, towards Tsunamis generation with the multi-layer solver

Fanshuo Ma, Mathilde Tavares, Pierre-Yves Lagrée, Stéphane Popinet

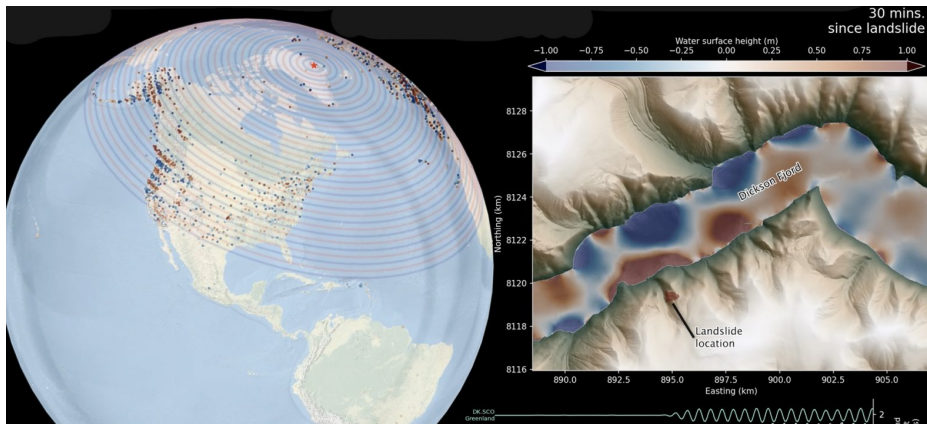
Institut Jean Le Rond d'Alembert, Sorbonne Université, France



Gravity flows tsunamis

Massive landslide in Groenland - september 2023

25 millions m^3 of rocks and ice mixture collapse into the Dickson Fjord $\Rightarrow$ 200 m wave

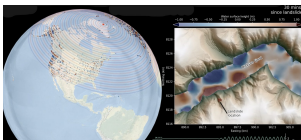


"Un énorme glissement de terrain déclenché par le changement climatique a fait vibrer la Terre pendant 9 jours" (IPGP - Institut de Physique du Globe de Paris)

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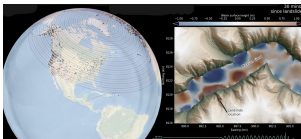
IPGP reminds us “the importance of better **anticipating** these events to **protect** populations and ecosystems”

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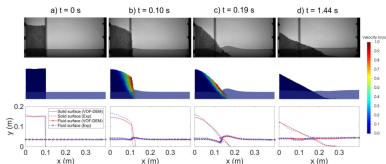


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From field to laboratory: geophysical flow approximations

- Shallow water approximation ($L \gg h$)
- 2D continuous model and /or discrete, often local

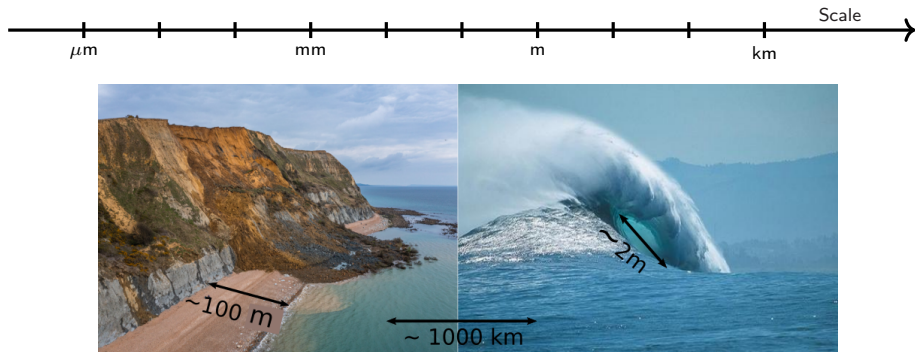


Granular collapse in water, expérience - top (Cabrera et al, JGR, 2020), numerical simulation - bottom (Nguyen, PoF, 2022)

What mechanism drives the formation of a wave when a **large amount of material** * enters the water

* Granular medium: glacier, rocks, volcano...

Gravity flow tsunamis – a multiscale and multiphase problem



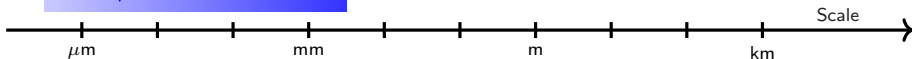
Landslide in the Dorset, UK (left), (*James Loveridge Photography*), wave breaking (right), (*google image*)

The problem description

- ① Granular material with complex rheology
- ② Wave formation and breaking
- ③ Imbibition, phase mixture, triple line front evolution...

Gravity flow tsunamis – a multiscale and multiphase problem

DEM/Navier-stokes

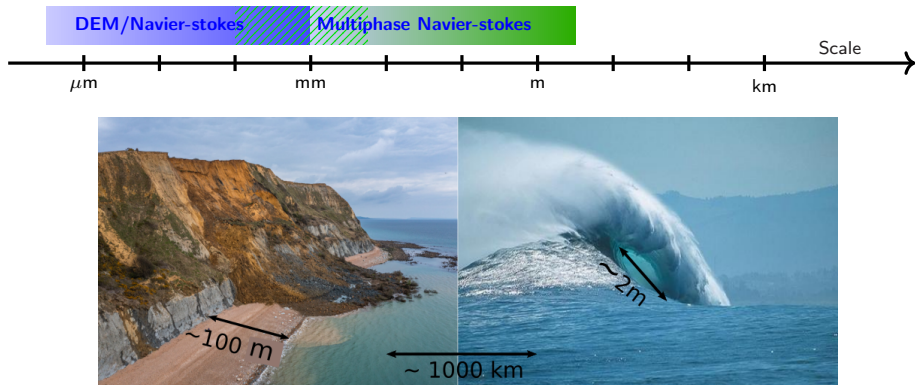


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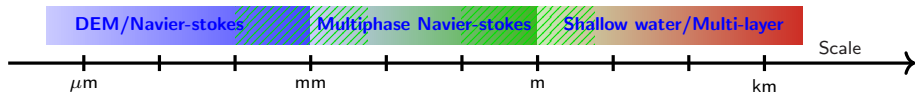


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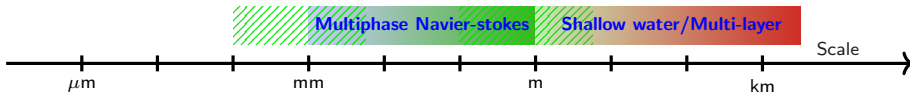
Gravity flow tsunamis – a multiscale and multiphase problem



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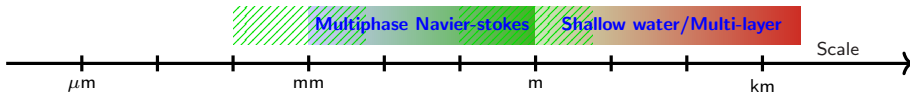
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Wave breaking

Multiphase Navier-Stokes

- ✗ Grid and refinement
- ✗ Complex numerical schemes
- ✗ High computational cost



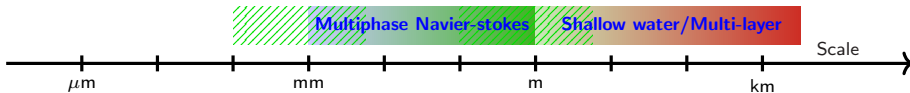
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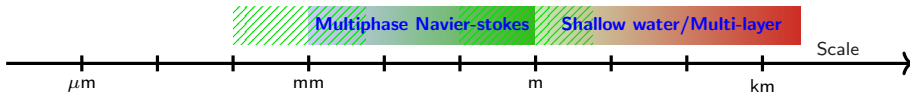
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Multi-layer model

- ✓ Fixed grid (almost)
- ✓ Simple numerical schemes
- ✓ Lower computational cost



Landslide in Dorset, UK



Wave breaking

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Multi-layer model

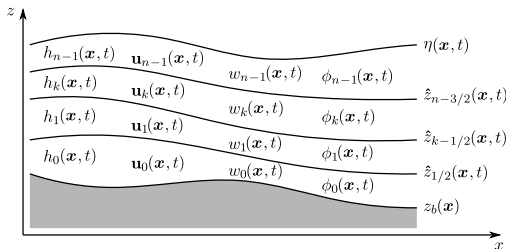
- ✓ Fixed grid (almost)
- ✓ Simple numerical schemes
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Proposition

Multi-layer solver for gravity flows and/or tsunamis formation



The great wave off Kanagawa (Wikipedia)



Definition of N layers (credit: <http://basilisk.fr/>)

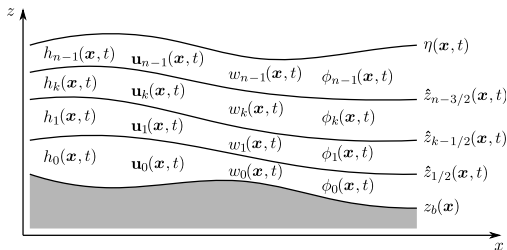
$$\begin{aligned}
 \partial_t h_k + \bar{\nabla} \cdot (h \bar{u})_k &= 0 \\
 \partial_t (h \bar{u})_k + \bar{\nabla} \cdot (h \bar{u} \bar{u})_k &= -g h_k \bar{\nabla}(\eta) - \bar{\nabla}(\Phi)_k + [\Phi \bar{\nabla}]_k \\
 \partial_t (h w)_k + \bar{\nabla} \cdot (h w \bar{u})_k &= -[\Phi]_k \\
 \bar{\nabla} \cdot (h \bar{u})_k + [w - \bar{u} \cdot \bar{\nabla} \hat{z}]_k &= 0
 \end{aligned}$$

With k the index of the layer, h_k its thickness, u_k , w_k , the horizontal and vertical velocity, g the gravity and Φ_k the non-hydrostatic pressure,

$$\eta \equiv z_b + \sum_k h_k \text{ and } \hat{z}_{k+1/2} \equiv z_b + \sum_{l=0}^k h_l$$



Landslide in Dorset, UK



Definition of N layers (credit: <http://basilisk.fr/>)

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With k the index of the layer, h_k its thickness, u_k , w_k , the horizontal and vertical velocity, g the gravity and Φ_k the non-hydrostatic pressure, ν the kinematic viscosity

$$\eta \equiv z_b + \sum_k h_k \text{ and } \hat{z}_{k+1/2} \equiv z_b + \sum_{l=0}^k h_l$$

Implicit viscous friction term:

$$\frac{(hu_I)^{n+1} - (hu_I)^*}{\Delta t} = \nu \left(\frac{u_{I+1} - u_I}{h_{I+1/2}} - \frac{u_I - u_{I-1}}{h_{I-1/2}} \right)^{n+1}$$

Problem

Complex rheology for granular flow $\Rightarrow$ viscosity $\nu \neq \text{cste}$

Neumann boundary conditions on the free surface

Navier slip on the bottom

$$\partial_z u|_t = \dot{u}_t$$

Default free slip $\dot{u}_t = 0$

$$u|_b = u_b + \lambda_b \partial_z u|_b$$

and no slip $u_b, \lambda_b = 0$

These equations can be expressed as the linear system

$$\mathbf{M}\mathbf{s}^{n+1} = \text{rhs}$$

where $\mathbf{M}$ is a tridiagonal matrix.

Implicit viscous friction term:

$$\frac{(hs_I)^{n+1} - (hs_I)^*}{\Delta t} = \left(\nu_{eq\,I+1/2} \frac{s_{I+1} - s_I}{h_{I+1/2}} - \nu_{eq\,I-1/2} \frac{s_I - s_{I-1}}{h_{I-1/2}} \right)^{n+1}$$

Development

Complex rheology for granular flow $\Rightarrow$ variable viscosity ν_{eq}

Neumann boundary conditions on the free surface

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$\mu(I)$ rheology (*GRD midi, Lagrée et al, JFM, 2011*) $\Rightarrow$ tangential and normal stress are linked through:
$$\tau = \mu(I)p$$

with the friction law

$$\mu = \mu_s + \frac{\Delta\mu}{1 + I_0/I} \quad \text{where} \quad I = d_g \frac{\partial u / \partial z}{\sqrt{\rho/\rho}}$$

I is the inertial number and d_g the grain diameter;
Coefficients μ_0 , $\Delta\mu$ and I_0 depend on the nature of the granular media

Equivalent rheology in the multi-layer model

The solver is tuned to consider non newtonian viscosity as a function of the shear

$$\nu_{eq_k} = \nu + \frac{\tau_z / \rho}{\left| \frac{\partial u}{\partial z} \right|_k}$$

Typically, ν_{eq_k} and $\left| \frac{\partial u}{\partial z} \right|_k$ also depends on the index layer k ;

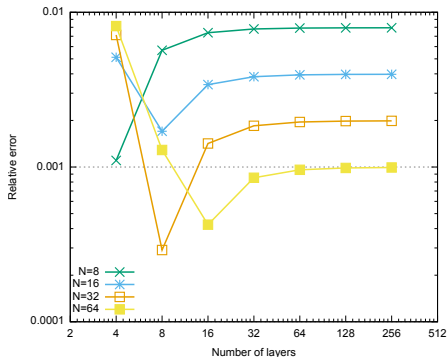
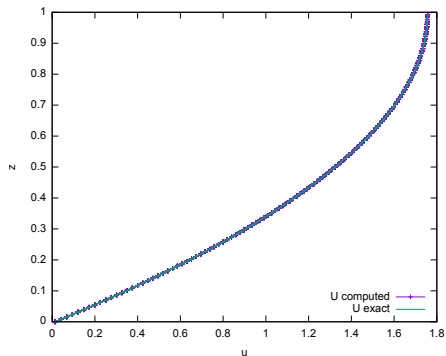
A shear regularisation is required to avoid infinite viscosity at rest

Periodic free surface with Newtonian flow

Poiseuille

A theoretical solution for the velocity:

$$u = \frac{G \sin \theta}{2\nu} (2 - z)z$$



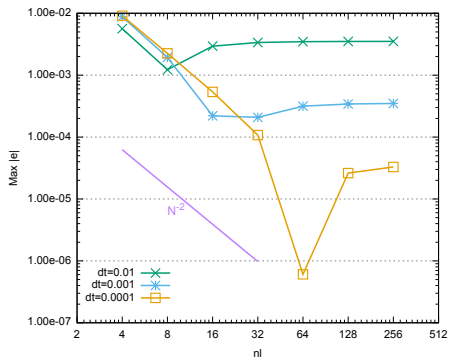
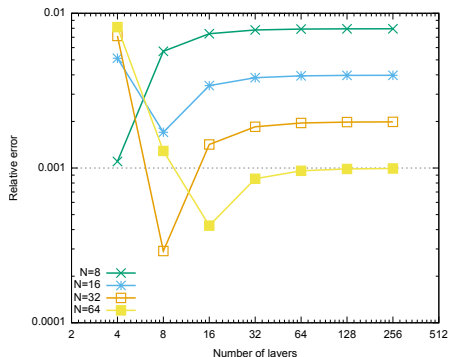
Comparison between analytical solution and numerical solution (l); order of convergence (r) - $N_x = 16$, $n_l = 128$, **Poiseuille**

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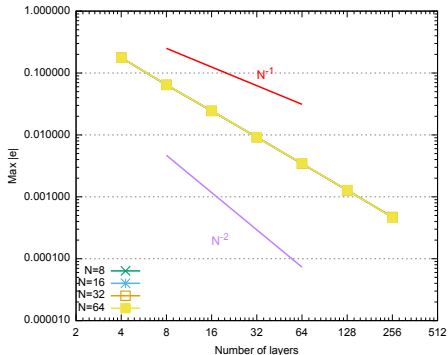
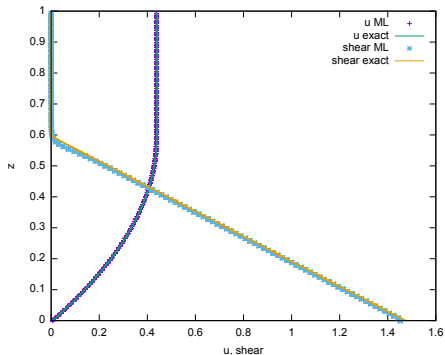
Order of convergence, influence of dt - Poiseuille

Periodic free surface with non Newtonian flow

Bingham

A theoretical solution for the velocity:

$$u = U \left(1 - \left(\frac{z - Z}{Z} \right)^2 \right) \quad \text{with} \quad U = \frac{(\tau_z - h(-G \sin \theta))^2}{2\mu(G \sin \theta)} \quad \text{and} \quad Z = h - \tau_z/(-G \sin \theta)$$



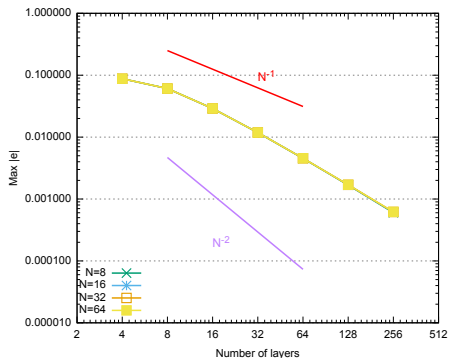
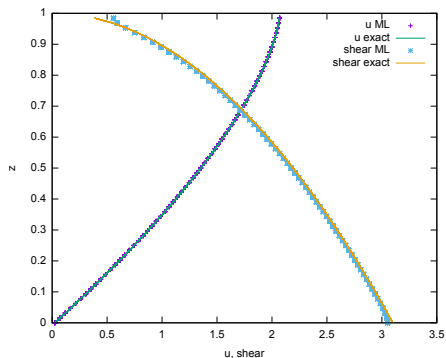
Comparison between analytical solution and numerical solution (l); order of convergence (r) - $N_x = 16$, $n_l = 128$, **Bingham**

Periodic free surface with non Newtonian flow

Bagnold

A theoretical solution for the velocity:

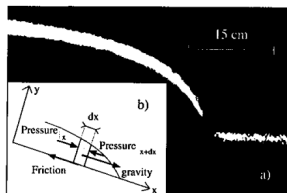
$$u = \frac{2}{3} \frac{l_\theta}{d_g} \sqrt{gh^3 \cos(\theta)} \left(1 - \left(1 - \frac{z}{h}\right)^{3/2}\right) \quad \text{with } l_\theta = \frac{l_0(\tan \theta - \mu_0)}{\mu_0 + \Delta\mu - \tan \theta};$$



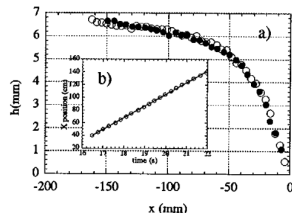
Comparison between analytical solution and numerical solution (I); order of convergence (r) - $N_x = 16$, $nl = 128$, **Bagnold**

Shape of a granular front down a rough inclined plane

An avalanche flowing along a mild slope



(a) Front illuminated by the laser sheet $\theta = 21^\circ$, $h_\infty = 9.5\text{mm}$,
(b) Forces on a elementary material slice



(a) Front profiles 50cm from the outlet $\circ$ and 150cm down the slope $\bullet$, (b)
Position of the front with time showing the constant propagation velocity.

A theoretical solution for the front shape (G. Saingier et al 2016)

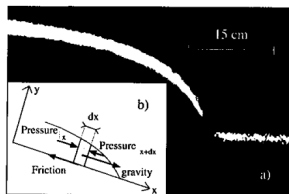
$$X = H - \frac{\frac{2}{3} \log(1 - \sqrt{H}) - \frac{1}{3} \log(H + \sqrt{H} + 1) + \frac{2}{\sqrt{3}} \arctan\left(\frac{2\sqrt{H}+1}{\sqrt{3}}\right) - (\alpha - 1)Fr^2 \ln\left(\frac{H^d}{(1-H^{3/2})^{2/3}}\right)}{d - 1}$$

With $Fr = \frac{u_0}{\sqrt{gh_0 \cos \theta}}$, u_0 the mean velocity; $H = h/h_0$; $X = x(\tan \theta - \mu_0)/h_0$;

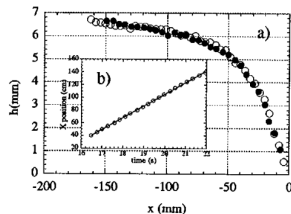
$$d = (\tan \theta - \mu_0)/\Delta\mu.$$

Shape of a granular front down a rough inclined plane

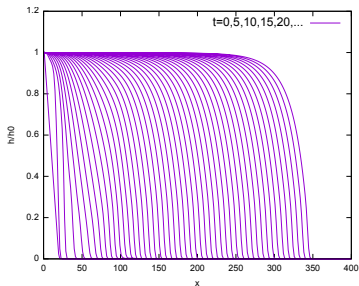
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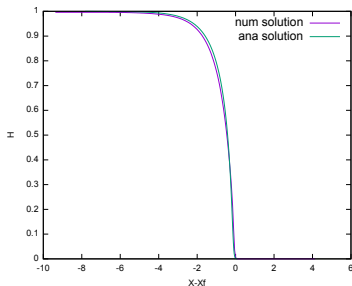
(a) Front illuminated by the laser sheet $\theta = 21^\circ$, $h_\infty = 9.5\text{mm}$,
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(a) Front profiles 50cm from the outlet $\circ$ and 150cm down the slope $\bullet$, (b)
Position of the front with time showing the constant propagation velocity.



Shape of the front as a function of time (l), comparison between analytical solution and numerical solution (r), $N_x = 256$, $n_l = 128$



Quick overview on multi-layer model for Gravity Flows

Ongoing works

- Variable viscosities in multi-layer solver
- Different existing rheologies
- Validation with several test cases and examples

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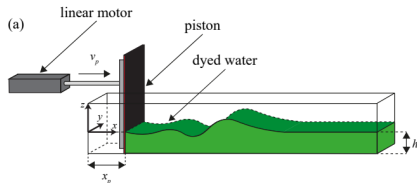
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- Variable viscosities in multi-layer solver
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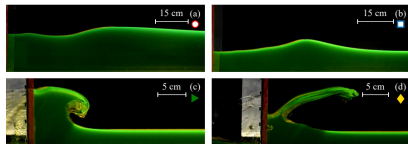
To do

- Time convergence of diffusion
- Extend Neumann to Navier BC on the free surface
- Two-phase problem (Poiseuille/Bingham) with variable viscosities in the Boussinesq approximation

Waves impulsed by granular collapse



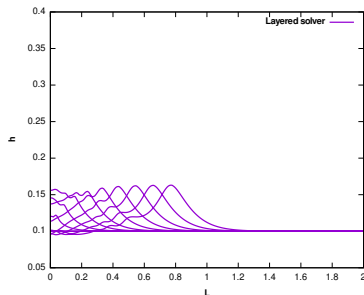
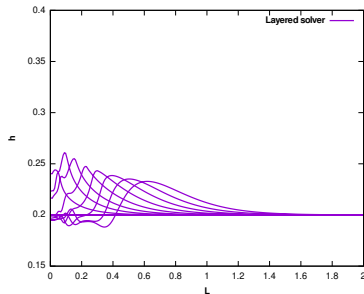
Experimental set-up, (Sarlin et al, JFM, 2025)



(a) A dispersive wave, (b) A solitary-like wave (Sarlin et al, JFM, 2025)

(a) $L = 7\text{cm}$, $U = 0$, 42m.s , $h = 20\text{cm}$, $Fr_p = 0,3$

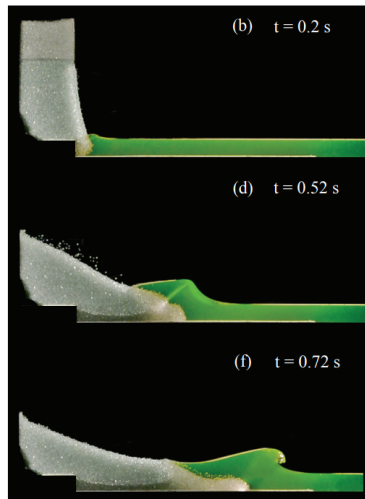
(b) $L = 15\text{cm}$, $U = 1,09\text{m.s}$, $h = 10\text{cm}$, $Fr_p = 0,47$



Gravity flows tsunamis

Physical challenge

Numerical features



Wave generation by granular collapse, $h_0 = 5\text{cm}$, (Robbe-saule et al, *JFM*, 2021)

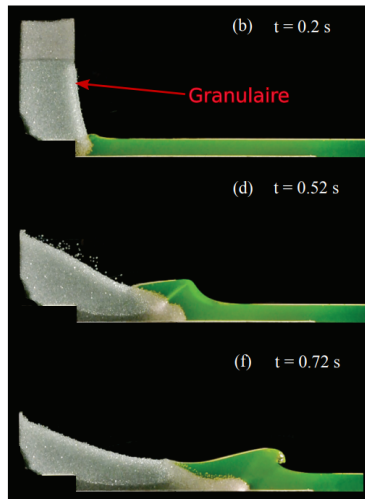
Gravity flows tsunamis

Physical challenge

- ❗ Granular rheology

Numerical features

- ▶ Multi-layer and variable viscosity



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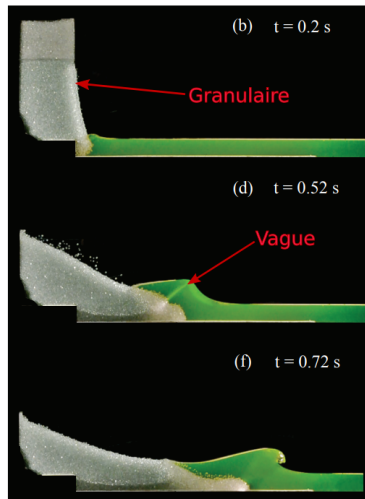
Gravity flows tsunamis

Physical challenge

- 1 Granular rheology
- 2 Wave formation

Numerical features

- ▶ Multi-layer and variable viscosity
- ▶ Multi-layer and **variable density**



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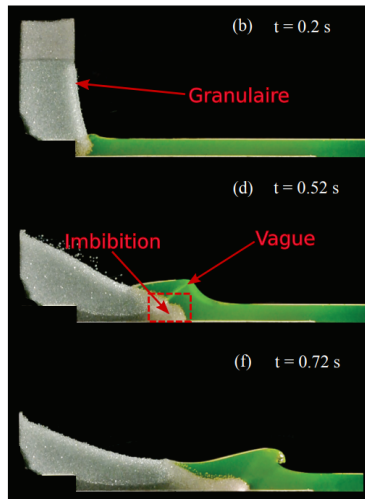
Gravity flows tsunamis

Physical challenge

- ❶ Granular rheology
- ❷ Wave formation
- ❸ Imbibition/porous medium

Numerical features

- ▶ Multi-layer and variable viscosity
- ▶ Multi-layer and **variable density**
- ▶ Triple line modeling



Génération de vague par effondrement d'une colonne granulaire,
 $h_0 = 5\text{cm}$, (Robbe-saule et al, JFM, 2021)

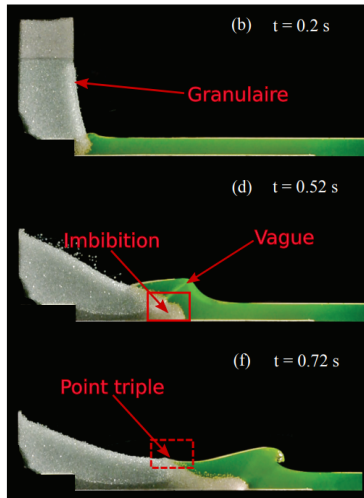
Gravity flows tsunamis

Physical challenge

- ❶ Granular rheology
- ❷ Wave formation
- ❸ Imbibition/porous medium
- ❹ Triple line

Numerical features

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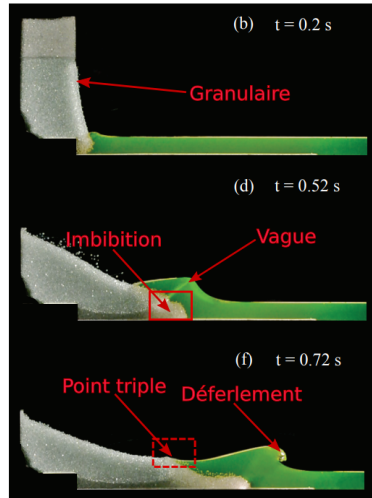
Gravity flows tsunamis

Physical challenge

- ❶ Granular rheology
- ❷ Wave formation
- ❸ Imbibition/porous medium
- ❹ Triple line
- ❺ Breaking wave

Numerical features

- ▶ Multi-layer and variable viscosity
- ▶ Multi-layer and variable density
- ▶ Triple line modeling
- ▶ Multi-layer and Navier-Stokes/Lagrangian particles



Génération de vague par effondrement d'une colonne granulaire,
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Thank you for your attention

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