

Impact-Induced Fragmentation of Liquid Rims

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Introduction: Ocean Sprays



Major pathways of ocean spray generation

Enhancing air-sea mass, momentum and energy transfer Regulation by breaking waves

• Significance of Ocean Sprays

- Pathways of ocean spray generation
 - Bursting of entrained bubbles Spume drop ejection under high winds
 - Splashing of plunging waves

[1] Veron, F. (2015). Ocean spray. Annual Review of Fluid Mechanics, 47, 507-538.
[2] Deike, L. (2022). Mass transfer at the ocean-atmosphere interface: The role of wave breaking, droplets, and bubbles. Annual Review of Fluid Mechanics, 54, 191-224.

Introduction: Splashing

Multiple splashing events [2]:

After impact ('forward splashing' [1]) Small number of drops

Secondary splashing ('backward splashing' [1])

Larger number of drops;

Small and vertically-projected drops Curbed at large surface tension

Generation mechanism unclear

Small-scale splashing & bubble bursting Last longer and produce larger drops





The secondary splashing mechanism

[1] Kiger, K. T., & Duncan, J. H. (2012). Air-entrainment mechanisms in plunging jets and breaking waves. Annual Review of Fluid Mechanics, 44, 563-596.

[2] Mostert, W., Popinet, S., & Deike, L. (2022). High-resolution direct simulation of deep water breaking waves: transition to turbulence, bubbles and droplets production. *Journal of Fluid Mechanics*, *942*, A27.

[3] Erinin, M. A., Wang, S. D., Liu, X., Liu, C., & Duncan, J. H. (2022). Spray Generation by Plunging Breakers--Part 2. Droplet Characteristics. arXiv preprint arXiv:2210.01923.

Introduction: Liquid Rim Collision





- Rim collision as a model for secondary splashing
- Experiments of Neel *et al.* (2020) [1]: Collision of rims travelling on liquid films $We_c = 66$: RT destabilization of lamellae Skewed size distribution functions $We < We_c$: Gamma distribution
 - $We > We_c$: Bessel-based distribution.



Experimental photographs of Neel et al. (2020)

Problem Configuration

Controlling Parameters

$$We \equiv \frac{\rho_g (2U_0)^2 d_0}{\sigma}, \qquad Oh \equiv \frac{\mu_l}{\sqrt{\rho_l d_0 \sigma}} = 0.01, \\\rho^* \equiv \frac{\rho_l}{\rho_g} = 833, \qquad \mu^* \equiv \frac{\mu_l}{\mu_g} = 55,$$

With perturbation:

$$\varepsilon_0 \equiv \frac{2\epsilon_0}{d_0}, \qquad N_{\max}$$

Interfacial Perturbation [1]

1) White noise series

2) Filter and keep only the lowest $N_{\rm max}$ modes

3) Re-normalize to a variation of ε_0

4) Two cylinders perturbed by different realisations with the same ε_0 and N_{max} .

Basilisk, Two-Phase NS Equation w. AMR

[1] Pal, S., Crialesi-Esposito, M., Fuster, D., & Zaleski, S. (2021). Statistics of drops generated from ensembles of randomly corrugated ligaments. *arXiv preprint arXiv:2106.16192*.



Configuration of the rim collision problem



Power density of the perturbation series



Self-similar evolution of sheet velocity and outer contour :

$$\begin{cases} \boldsymbol{u}_{\boldsymbol{y}} = \frac{\boldsymbol{y}}{t} \\ \frac{ht}{R_0 \tau} = f\left(\frac{\boldsymbol{y}}{U_0 t}\right) \end{cases}$$

Independent of impact *We* and perturbation waveform

Implication: sheet expansion mostly inviscid, fluid parcel travels at constant velocity [1]



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Vertical velocity within liquid sheets

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[1] Wang, Y., & Bourouiba, L. (2017). Drop impact on small surfaces: thickness and velocity profiles of the expanding sheet in the air. Journal of Fluid Mechanics, 814, 510-534.

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Rim Kinematics



Scaling model for rim position and thickness:





Vertical position (left) and thickness (right) of the rim

Ligament Growth







 $-We = 160, t/\tau_{cap} = 0.36$

- 160, 0.54

Vertical velocity within sheets and ligaments

Lai, C. Y., Eggers, J., & Deike, L. (2018). Bubble bursting: Universal cavity and jet profiles. *Physical review letters*, 121(14), 144501.
 Gekle, S., & Gordillo, J. M. (2010). Generation and breakup of Worthington jets after cavity collapse. Part 1. Jet formation. Journal of fluid mechanics, 663, 293-330.



Ligaments merging on the corrugated rim

0.65



• Wang & Bourouiba (2018):

Ligament merging observed where the inflow is not perpendicular to the rim

Driven by the tangential component of inflow velocity

 $u_{\rm drift} = (u(R,t) - \dot{R})\sin\theta$



Ligaments merging on the corrugated rim

0.55

x/L

0.6

0.65



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 $\begin{array}{c} 0.12 \\ 0.08 \\ 0.04 \\ 0 \\ 0.5 \\ 0.5 \\ x/L \end{array}$

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Ligaments merging on the corrugated rim

$$\frac{1}{N_{\rm lig}^2} = AWe^{0.25}t + B$$



Ligament number N_{lig} at different N_{max} (left) and We (right)

Primary Drop Statistics



Total number (left) and diameter (right) of primary fragments

Primary Drop Statistics



End-pinching of Worthington jets: $\langle d/w \rangle \approx 1.5$ [1][2]



Wang, Y., & Bourouiba, L. (2018). Unsteady sheet fragmentation: droplet sizes and speeds. *Journal of Fluid Mechanics*, 848, 946-967.
 Gordillo, J. M., & Gekle, S. (2010). Generation and breakup of Worthington jets after cavity collapse. Part 2. Tip breakup of stretched jets. *Journal of fluid mechanics*, 663, 331-346.









Conclusions and Future Work

- ✓ 3D simulations of cylinder collision;
- ✓ Liquid sheet expansion:
 Self-similar models for u_y and h
- ✓ Rim Kinematics:

Scaling models for y_{rim} and b_{rim}

✓ Ligament Dynamics:

Linear growth and ballistic region

Theoretical model for ligament merging

✓ Fragment statistics: agreement with existing results for large drops

- Refine the ligament merging model;
- Establish grid convergence for fragment statistics;
- Explore connection with the drop size distribution of wave splashing.

Thanks for your attention!