Falling liquid film control via linear quadratic regulation

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Falling liquid film control

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Introduction Motivation



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Introduction Motivation



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Introduction Motivation



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Thin fluid film falling down an inclined plane



Thin fluid film falling down an inclined plane

Our aim is to stabilise the uniform film solution by injecting and removing fluid from the base at a finite number of actuators.



Thin fluid film falling down an inclined plane

Navier-Stokes flow in the fluid

$$\begin{aligned} ℜ(u_t + uu_x + vu_y) = -p_x + 2 + u_{xx} + u_{yy}, \\ ℜ(v_t + uv_x + vv_y) = -p_y - 2\cot\theta + v_{xx} + v_{yy}, \\ &u_x + v_y = 0. \end{aligned}$$

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Boundary conditions at the base

$$u=0, \quad v=f(x,t).$$

Introduction

Navier-Stokes film



Thin fluid film falling down an inclined plane

At the interface, y = h(x, t), the nonlinear dynamic stress balance

$$(v_x + u_y)(1 - h_x^2) + 2h_x(v_y - u_x) = 0,$$

 $p - \frac{2}{1 + h_x^2}(v_y + u_x h_x^2 - h_x(v_x + u_y)) = -\frac{1}{Ca} \frac{h_{xx}}{(1 + h_x^2)^{3/2}},$

and the kinematic boundary condition

$$h_t = v - uh_x$$
.

Introduction



- multi-phase flow
- complex boundary conditions
- highly nonlinear
- computationally expensive

The most general feedback control problem looks like

$$x_t = \mathcal{A}x + \mathcal{B}u, \qquad u = \mathcal{K}y, \qquad y = \mathcal{C}x.$$

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Introduction Linear quadratic regulator control

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We now need to choose K so that A + BK has no positive eigenvalues. The choice of K is currently not unique, so we introduce a quadratic cost

$$c = \int_0^\infty x^\mathsf{T} U x + u^\mathsf{T} V u \, \mathrm{d}t,$$

thus forming an LQR problem.

Introduction Problem of control



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LQR controls

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Reduced order model

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$$\begin{aligned} h_t + q_x &= f, \\ \frac{2Re}{5}h^2q_t + q &= \frac{h^3}{3}\left(2 - 2h_x\cot\theta + \frac{h_{xxx}}{Ca}\right) \\ &+ Re\left(\frac{18q^2h_x}{35} - \frac{34hqq_x}{35} + \frac{hqf}{5}\right). \end{aligned}$$

These are the weighted-residual integral boundary layer equations.

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$$+ Re\left(\frac{18q^2h_x}{35} - \frac{34hqq_x}{35} + \frac{hqf}{5}\right).$$

These are the WR equations.

Reduced order model



Development of travelling wave for **Navier-Stokes** and **WR** systems. Re = 10, Ca = 0.05.

Linearisation

$$\begin{aligned} h_t + q_x &= f, \\ \frac{2Re}{5}h^2q_t + q &= \frac{h^3}{3}\left(2 - 2h_x\cot\theta + \frac{h_{xxx}}{Ca}\right) \\ &+ Re\left(\frac{18q^2h_x}{35} - \frac{34hqq_x}{35} + \frac{hqf}{5}\right). \end{aligned}$$

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Linearisation

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$$+ Re\left(\frac{18q^2h_x}{35} - \frac{34hqq_x}{35} + \frac{hqf}{5}\right).$$

These equations are still very nonlinear. Assuming that any perturbations from the uniform film are small

$$h = 1 + \delta \hat{h}, \qquad q = \frac{2}{3} + \delta \hat{q}, \qquad f = \delta \hat{f},$$

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$$h = 1 + \delta \hat{h}, \qquad q = \frac{2}{3} + \delta \hat{q}, \qquad f = \delta \hat{f},$$

we have

$$\hat{h}_t = -\hat{q}_x + \hat{f},$$

$$\hat{q}_t = \left[\frac{5}{Re} + \left(\frac{4}{7} - \frac{5\cot\theta}{3Re}\right)\partial_x + \frac{5}{6ReCa}\partial_{xxx}\right]\hat{h} - \left[\frac{5}{2Re} + \frac{34}{21}\partial_x\right]\hat{q} + \frac{1}{3}\hat{f}.$$

Discretisation

Finally we can discretise

$$\begin{bmatrix} h \\ q \end{bmatrix}_t = \begin{bmatrix} J_{hh} & J_{hq} \\ J_{qh} & J_{qq} \end{bmatrix} \begin{bmatrix} h \\ q \end{bmatrix} + \begin{bmatrix} \Psi_h \\ \Psi_q \end{bmatrix} f,$$

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Or, in more concise notation,

$$\xi_t = (J + \Psi K) \xi$$

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Numerical experiments

What does this actually look like in practice?

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Numerical experiments

Gain matrix



Feedback gain for a single **actuator**. *Re* various, Ca = 0.05.

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Initial development of a travelling wave. Re = 15, Ca = 0.05.

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Measurement of the **height**. Re = 15, Ca = 0.05.

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Image: A match a ma



Computation of **controls**. Re = 15, Ca = 0.05.

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Image: A match a ma

Controls stabilising the uniform film. Re = 15, Ca = 0.05.

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Controls attempting to stabilise the uniform film. *Re* various, Ca = 0.05.

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What predictions can we make about the stabilisability of the system?

We can't make any predictions about the stabilisability of the full Navier-Stokes system.

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Shifting to Fourier space and explicitly separating the unstable modes

$$\tilde{\xi}_t = \begin{bmatrix} \tilde{J}_u + \tilde{\Psi}_u \tilde{K}_u & 0\\ \tilde{\Psi}_s \tilde{K}_u & \tilde{J}_s \end{bmatrix} \tilde{\xi}.$$

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So the number of controls should exceed the number of unstable modes.

Stability analysis Unstable modes

If we look at a unimodal perturbation $h=1+\hat{h}e^{ikx+\lambda t}$ we have

$$\lambda^2 + \left(\frac{5}{2Re} + \frac{34}{21}ik\right)\lambda + \left(\frac{5}{Re}ik - \left[\frac{4}{7} - \frac{5\cot\theta}{3Re}\right]k^2 + \frac{5}{6ReCa}k^4\right) = 0.$$

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Rescaling to allow for $L \neq 2\pi$ we can compute the unstable mode count

$$n_{\rm u} = 1 + 2 \left\lfloor \frac{L}{2\pi} \sqrt{Ca \left(\frac{8}{5}Re - 2\cot\theta\right)} \right\rfloor$$

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Image: A matrix

Numerical comparison



In practice the controls outperform the linear predictions.

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• Optimal feedback control for complex systems is achievable

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- Optimal feedback control for complex systems is achievable
- Controls function well outside the range of model validity

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- Controls exceed expected performance

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Next steps:

• 3D flows

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- 3D flows
- Alternative actuator mechanisms

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- 3D flows
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- Restricted observations

- Optimal feedback control for complex systems is achievable
- Controls function well outside the range of model validity
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- 3D flows
- Alternative actuator mechanisms
- Restricted observations
- Infinite-dimensional control

More detail

Preprint available on arXiv



Code available on GitHub



arxiv.org/pdf/2301.11379

github.com/OaHolroyd/falling-filmcontrol/tree/paper-dec-2022

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